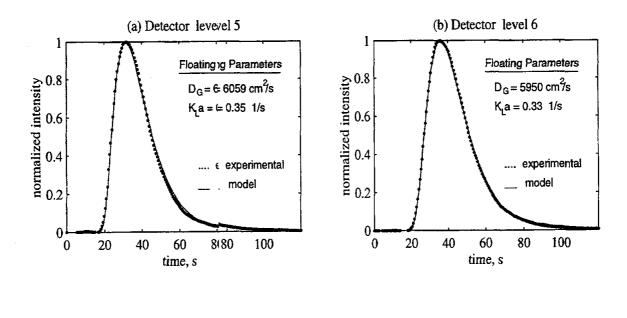
Table 8: Case 2b: Parameterers with Fixed H for Gas Phase Tracer Experiments with NDG Measurements for Gas 3 Holdup at Upper Detector Levels

Run No.	$U_{\ell}J_{G}$	Gas	D_L	Det.		Parameters	
		Holdup		Lev.	H	$\overline{D_G}$	K_{La}
	cmn/s		cm^2/s			cm^2/s	s^{-1}
14.6-1	25.5.3	0.42	4696	5	5.86	6059	0.35
		0.45		6		5950	0.33
		0.50		7		5504	0.13
14.6-2		0.42		5	5.86	5805	0.71
		0.45		6		5755	0.47
		0.50		7		4825	0.15
14.7-3	14.4.3	0.37	3052	5	5.86	2416	0.22
1		0.39		6		2472	0.21
		0.41		7		1916	0.14
14.7-4		0.37		5	5.86	2403	0.21
		0.39		6		2399	0.17
	1	0.41	İ	7		2306	0.23
14.8-5	36.6.0	0.41	5925	5	8.11	9621	2.63
		0.44		6		8325	1.54
		0.46		7		6720	0.832
14.8-6		0.41		5	8.11	9088	3.32
		0.44		6		7710	1.85
		0.46		7		6970	0.84



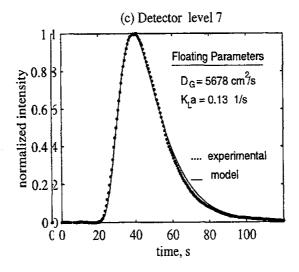


Figure 9: Gas Phase IrImpulse Response for Run 14.6-1, Injection Time 12.6s

5.2.3 Case 3: Model with Three Floating Parameters: ϵ_G , D_G and $K_L a$

In order to find out what the model predicts as average gas holdup, the data are fitted with the other choice of three parameters by floating ϵ_G , D_G and $K_L a$, and fixing H at its thermodynamic value. As one would expect the fits are very good. The values for the obtained parameters are shown in Table 9. The estimated values of ϵ_G are higher than those nineasured by DP. The discrepancy is larger in the upper part of the column than at the lower level, but exists for the lower portions of the column as well. For some cases in runs 14.7 and 14.8, the estimated holdup values are larger than both DP and NDG measured values! Therefore, little can be said about the D_G and $K_L a$ values from this case of parameter estimation, unless we can independently confirm a proponounced axial solids concentration profile which would have affected the measured ggas holdup profiles.

5.2.4 Case 4: Model with Three Floating Parameters: D_G , $K_L a$ and α

In the analysis done so fair for the gas phase tracer experiments, the tracer concentration $C_t(t,z)$ was assummed to be given by Equation 7. This relationship holds if one assumes no cross-sectitional variation of tracer and gas holdup (as considered by the one-dimensional modelel). However, in reality there is a radial variation of the gas holdup profile (Hills, 19771; Kumar et al., 1994). This affects the measurement of the total argon tracer concentration because the detector gets its major source of radiation from the tracer closse to the wall. If the phase holdups and concentrations C_G and C_L are uniform through out the cross section, then the total concentration of tracer is the average of the gas and liquid phase concentrations, weighted by the respective holdups. If, however, the distribution of phases is not cross-sectionally uniform, but has more liquid at t the wall, as is the actual case, then the average should not be weighted by the phase e holdups, but must also account for the relatively higher contribution of the liquid, duue to higher volume fraction of the liquid at the wall. In other words $C'_t(t,z) = \alpha \epsilon_g C_{GG} + \epsilon_L C_L$ where $\alpha \leq 1.0$.

It is not possible to know v what α is without modeling the radiation received by the detector for a given spatial didistribution of the phases and tracer concentration. This in itself is a rigorous and timme-consuming procedure. For a preliminary assessment of the effect of an uneven crooss-sectional distribution, the model fitting is performed using α as the third floating v parameter, instead of e or v.

Table 9: Case 3: Parammeters for Gas Phase Tracer Experiments with Fixed ${\bf H}$

Run No.	\overline{U}_G		$\overline{D_L}$	Det.	Model	Paran	
		H		Lev.	D_G	ϵ_G	K_La
	c cm/s		cm^2/s		cm^2/s		s^{-1}
14.6-1	2 25.3	5.86	4696	1	20000	0.20	3.158
				2	6790	0.43	2.495
	:			3	7879	0.39	1.855
				4	5725	0.43	0.36
				5	7473	0.41	1.45
				6	6256	0.45	0.44
1			-	7	6364	0.49	0.67
14.6-2				1	15000	0.24	4.04
				2	5213	0.46	1.32
:			1	3	6361	0.41	2.19
				4	4317	0.48	0.22
				5	5937	0.43	0.77
				6	5574	0.47	0.33
				7	5246	0.49	0.43
14.7-3	1 14.3	5.86	3052	1	7340	0.34	0.006
		İ		2	4094	0.39	1.39
	, ,			3	4533	0.36	1.58
				4	3628	0.38	0.35
				5	2970	0.42	0.32
	1	j		6	2462	0.45	0.16
				7	2426	0.45	0.23
14.7-4				1	5929	0.35	0.006
				2	3329	0.38	0.28
				3	4170	0.34	0.42
1			Ì	4	2954	0.40	0.16
				5	2553	0.42	0.18
]		6	3111	0.44	0.13
		1		7	2606	0.46	0.14
14.8-5	3 36.0	8.11	5925	1	20000	0.20	1.80
		1	1	2	7235	0.46	0.20
			1	3	7963	0.43	0.29
)		4	7425	0.49	0.10
				5	7320	0.49	0.17
				6	7135	0.51	0.10
				7	6689	0.52	0.08
14.8-6				1	20000	0.22	1.58
				2	9442	0.39	0.28
				3	9800	0.38	0.34
				4	8518	0.46	0.09
				5	7027	0.49	0.12
				6	6580	0.51	0.09
				7	7079	0.54	0.15

Table 10: Case 4: PParameters for Gas Phase Tracer Experiments

Run No.	UU_G	Gas	D_L	Det.	Model	Param	eters
		Holdup		Lev.	D_{G}	$K_L a$	α
	cmn/s		cm^2/s		cm^2/s	s^{-1}	
14.6-1	25.5.3	0.39	4696	1	9800	0.08	1.0
		0.39		2	7153	2.28	0.198
		0.38		3	7580	2.62	1.00
		0.37		4	5458	0.40	0.22
		0.38		5	7409	0.82	0.17
		0.39		6	6607	0.34	0.07
		0.40		7	6795	0.25	0.00
14.6-2	***	0.39		1	9800	0.44	1.00
		0.39		2	4854	0.74	0.23
		0.38		3	6428	2.32	0.29
		0.37		4	3763	0.37	0.10
		0.38		5	5863	0.54	0.11
ļ		0.39		6	5729	0.29	0.07
		0.40		7	5544	0.23	0.00
14.7-3	14.4.3	0.34	3052	1	4949	0.64	1.00
1		0.33		2	4678	0.78	0.15
		0.32		3	5026	0.99	0.16
	}	0.32		4	3697	0.35	0.15
		0.33		5	3216	0.30	0.00
		0.34		6	2467	0.16	0.00
		0.35		7	2274	0.12	0.00
14.7-4		0.34		1	4283	0.51	1.00
		0.33		2	3113	0.29	0.30
		0.32]	3	4143	0.40	0.38
		0.32		4	2571	0.26	0.15
		0.33		5	2343	0.29	0.04
		0.34		6	2661	0.16	0.00
	<u> </u>	0.35		7	2499	0.12	0.00
14.8-5	36.6.0	0.38	5925	1	17470	0.14	1.00
		0.37		2	6007	0.43	0.20
		0.37		3	7216	0.44	0.25
		0.37		4	5135	0.28	0.12
		0.37		5	6025	0.30	0.06
		0.38		6	5671	0.24	0.04
		0.38		7	5863	0.22	0.02
14.8-6		0.38	T	1	18973	0.11	1.0
		0.37		2	8893	0.37	0.38
		0.37		3	9740	0.38	0.42
		0.37		4	5873	0.23	0.17
		0.37		5	4722	0.28	0.08
		0.38		6	4362	0.22	0.05
		0.38		7	7151	0.23	0.05

Again, the fits are good for all conditions. The parameters are reported in Table 10. There is a distinct thrend in the values of α , which is zero at the upper detector levels, and increases with a decrease in height. Such low values of α at the higher levels suggest that the detectors at these levels see only the liquid and no gas. These results can have meaning oonly if this can be verified.

In order to verify the vivalidity of the results in Case 4, the radiation detected by a detector is modeled by accounting for the cross-sectional variation of gas holdup in the column. Details of this s calculation are shown in Appendix II. Once an estimate of C_G and C_L is obtained from the model, it is used in the radiation model to calculate the total concentration C_{t-1} .

Figure 10 shows the ressult of fitting the model response to the experimental data for detector level 7 of Run 1 14.6-1. The fit is similar to the case shown in Figure 8g, for the case of two-parameteter fitting, using holdup estimates from DP measurements. This implies that the total tracer concentration C_t resulting from the radiation calculation is approximately the same as defined by Equation 7. The results of Case 4, with $\alpha \sim 0$, are therefore not valid. The reason for this is that the concentration of tracer in the liquid phasse is so low that the influence of high liquid holdup at the wall is offset by the low lidiquid tracer concentration (refer to Figure 11 and Figure A.2.2 in Appendix II). As seen in Figure A.2.2, the contribution of radiation at these low concentrations of liquidid tends to become more uniform across the cross-section, and hence, the effect of the assumed radial gas holdup profile on the results is not very significant. This justififies the use of Equation 7 for this specific gas tracer of low solubility.

So far we have considered a combination of D_G , $K_L a$, H, and ϵ_G as the floating parameters of the model. For Cases 1, 3 and 4, where three floating parameters are used, the fits are goodd, but the values of H, ϵ_G and α for the respective cases do not compare well with ϵ corresponding values from independent measurements (or methods of estimation). Four Case 2, when NDG-based values for the holdup are used, the fits are reasonably goodd with two floating parameters - D_G and $K_L a$. As discussed earlier, the poor fits obtainined when using ϵ_G from DP measurements are due to the differences in the mean residence times between the experimental data and model predicted results.

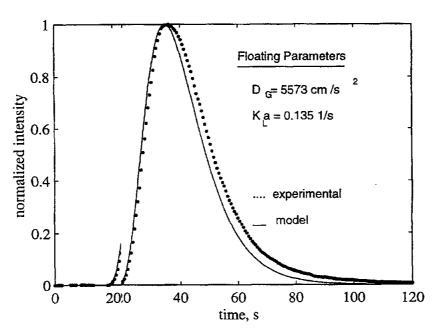


Figure 10: Gas Phase ImpulsIse Response for Run 14.6-1 at Detector Level 7, Injection Time 14.6s (Model I Uses Simulation of Incident Radiation)

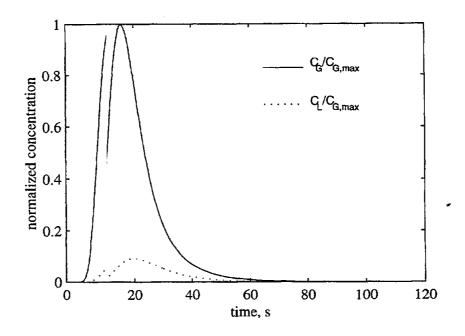


Figure 11: Model Calculated 1 Tracer Concentration in Gas and Liquid Phase (Run 14.6-1 Detector Level 7)

5.2.5 Case 5: Model with Changing Gas Flow Rate and two Floating Parameters, D_G and $K_L a$

Due to reaction in the slalurry bubble column reactor, there is a reduction in the gas volumetric flow rate. This 3 affects the superficial gas velocity along the length of the reactor. The conversions fefor the runs considered in this study result in an average reduction of around - 18%% in the gas flow rate. The axial dispersion model in its present form can not account for the axial variation of U_G . However, for the purpose of this analysis, the reducticion in gas flow rate is considered by assuming a linear change (decrease) in U_G with axialal position along the column, since the inlet and outlet gas flowrate are known from experimental data. This is just an approximation. Once the U_G at the intermediatete axial positions corresponding to the detector levels are evaluated, these values, $U_{G,G,z}$, are used in the model to predict the tracer distribution at different axial positions,;, and fit the experimental data. The average value of $U_{G,z}$ at a given axial position exam be used; however since the axial profile of $U_{G,z}$ is not known, the present variation in $U_{G,z}$ is considered. Thus, instead of using the inlet gas velocity, as done in Case 2,1, we now employ the gas velocity estimated at the detector level. The resulting parameters are reported in Table 11. Figure 12 shows the fits of the model to experimental data for Run 14.6 - 1. For this case, with only two floating parameters, the fits are much better than for Case 2 (using DP measurements for ϵ_G and inlet U_G), although they are not as perfect as for the three floating parameter cases. However, these fits aare acceptable, since only two floating parameters are used.

5.2.6 Case 6: Model with three Floating Parameters, D_G , $K_L a$ and U_G

As a final case, regressicion is performed to fit the experimental data to the model with three floating parameters, a position-dependent superficial gas velocity, $U_{G,z}$, D_G and $K_{L}a$. This is donne in order to estimate the average variation of U_G with axial position z. The resulting fits are as good as the other fits with three floating parameters. The values old the obtained parameters are shown in Table 12. The values of $U_{G,z}$, in general suggest a monotonic decrease of the gas velocity with axial position, especially at the higher detector levels, as expected.

Therefore this set of parameters for D_G , $K_L a$ and U_G , along with the independent estimates of ϵ_G and H, foorm the most reliable set of parameters. Although three floating parameters used, tithe values of $U_{G,z}$ obtained are reasonable, and fall within the range of experimentall values. This case is therefore the one that should be considered for analysis of tithe model parameters.

Table 11: Case 5: Parameters: for Gas Phase Tracer Experiments (Using Variable U_G)

Run No.	U_{CG}	Gas	D_L	Det.		Param	eters
		Holdup	-	Lev.	H	D_G	$K_L a$
	cmn/s	<i>T</i>	cm^2/s			cm^2/s	s^{-1}
14.6-1	24.8.82	0.39	4696	1	5.86	9025	0.006
	24.2.29	0.39		2		5310	0.63
	24.6.08	0.38		3		5577	0.28
1	23.3.15	0.37		4		4940	0.26
	22.2.29	0.38		5		5433	0.12
	21.1.9	0.39		6		4966	0.13
	21.1.5	0.40		7		4844	0.14
14.6-2	24.8.82	0.39		1	5.86	7599	0.006
1	24.2.29	0.39		2		4908	1.56
	24.0.08	0.38		3		4677	0.46
	23.3.15	0.37		4		3937	0.57
	22.2.29	0.38		5		4457	0.19
	21.1.9	0.39		6		4604	0.14
	21.1.5	0.40		7		4154	0.18
14.7-3	14.4.0	0.34	3052	1	5.86	5000	3.52
1	13.3.6	0.33		2	}	3410	0.78
}	13.3.4	0.32		3	Į l	2583	0.42
	12.2.7	0.32		4		3156	0.35
ŀ	12.2.2	0.32	ļ	5		3017	0.40
	11.1.8	0.33	ĺ	6		2990	0.51
	11.1.7	0.33		7		2100	0.28
14.7-4	14.4.0	0.33		1	5.86	5500	3.52
	13.3.6	0.33		2		3261	0.50
	13.3.4	0.32	1	3		2762	0.29
	12.2.7	0.32		4		2871	0.33
	12.2.2	0.32		5		2928	0.43
	11.1.8	0.33		6		2982	0.39
	11.1.7	0.33		7		2867	0.35
14.8-5	35.5.3	0.38	5925	1	8.11	13922	0.002
	34.4.4	0.37		2		8413	1.70
	34.4.0	0.37		3		7894	0.78
	32.2.5	0.37		4		6584	0.57
	31.1.0	0.37		5		6559	0.50
	30.0.1	0.38		6		5777	0.36
	29.9.6	0.38		7		5309	0.22
14.8-6	35.5.3	0.38		1	8.11	16533	0.002
	34.4.4	0.37		2		8691	0.27
[34.4.0	0.37		3		9119	0.25
	32.2.5	0.37		4		7325	0.28
	31.1.0	0.37		5		5792	0.40
	30.0.1	0.38		6		5108	0.29
	29.9.6	0.38		7	1	5765	0.30

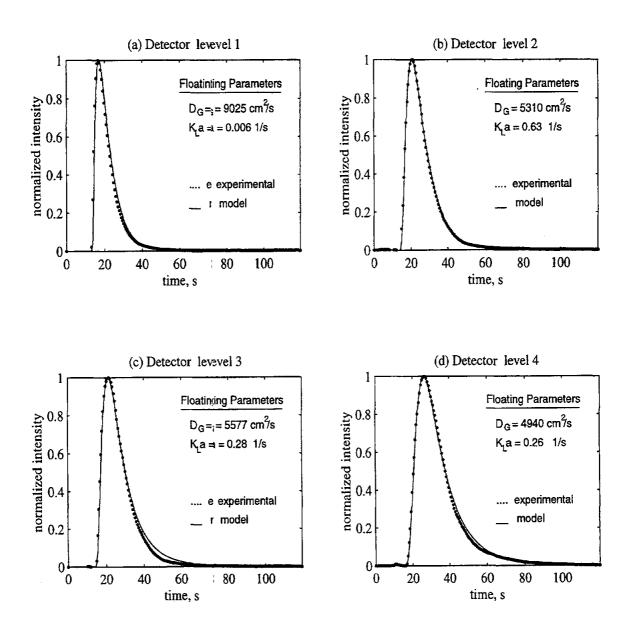
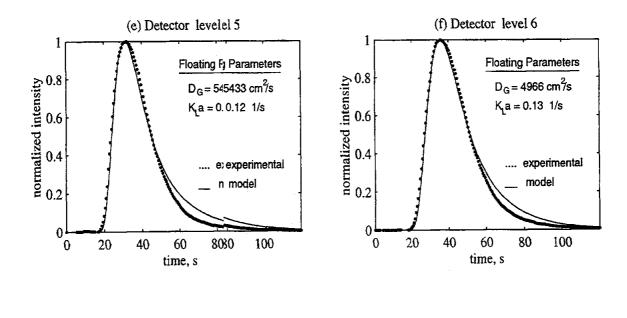


Figure 12: Gas Phase Impulse Response for Run 14.6-1 using Position Dependent Superfificial Gas Velocity, Injection Time 12.6s



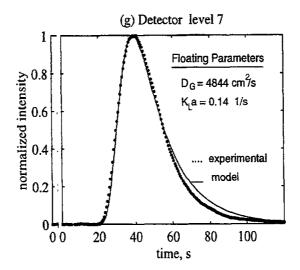


Figure 12 (contd.): Gas PPhase Impulse Response for Run 14.6-1 using Position Dependent Supperficial Gas Velocity, Injection Time 12.6s

Table 12: Case 6: : Parameters for Gas Phase Tracer Experiments

Run No.	U_G	Gas	D_L	Det.	Model	Parame	eters
		Holdup		Lev.	D_{G}	$K_L a$	U_G
	cicm/s		cm^2/s		cm^2/s	s^{-1}	cm/s
14.6-1	2 25.3	0.39	4696	1	9800	0.008	25.3
		0.39		2	6907	0.60	24.8
		0.38	{	3	6970	1.05	25.3
}		0.37		4	5282	0.45	23.9
		0.38		5	6974	0.36	24.4
}		0.39		6	5817	0.49	23.4
		0.40		7	5885	0.51	22.7
14.6-2		0.39		1	9800	0.001	25.3
		0.39		2	4536	0.99	24.1
		0.38		3	6239	0.63	25.0
		0.37		4	3717	0.44	22.8
		0.38		5	5500	0.23	23.8
j		0.39		6	5173	0.42	23.2
		0.40		7	4798	0.53	22.6
14.7-3	1 14.3	0.34	3052	1	4345	0.23	14.3
		0.33		2	3104	0.36	13.3
-		0.32		3	4082	0.44	13.9
		0.32		4	3449	0.43	13.2
	1	0.33		5	2714	0.47	12.4
		0.34		6	2133	0.33	12.1
		0.35		7	2157	0.43	12.0
14.7-4		0.34		1	3627	0.18	14.3
		0.33	1	2	3104	0.36	13.3
		0.32	}	3	4081	0.44	13.9
1		0.32	i i	4	2647	0.29	12.8
į		0.33		5	2230	0.34	12.4
		0.34		6	2286	0.34	12.1
		0.35		7	2288	0.31	11.9
14.8-5	3 36.0	0.38	5925	1	16150	0.003	36.0
		0.37		2	6250	0.51	32.7
		0.37		3	7205	0.54	33.3
	1	0.37		4	6127	0.21	30.9
		0.37		5	6226	0.45	30.6
		0.38		6	5722	0.33	30.0
		0.38		7	5348	0.30	29.9
14.8-6		0.38		1	19247	0.002	36.0
		0.37		2	9044	0.37	34.9
		0.37		3	9660	0.50	35.0
		0.37		4	7281	0.15	31.3
		0.37		5	5636	0.32	30.6
		0.38		6	5092	0.28	30.0
		0.38		7	5868	0.35	29.9

5.2.7 Discussion of Results

The parameters that truly need to be estimated from this model, are the gas phase dispersion coefficient LD_G and the gas-liquid mass transfer coefficient K_La . For the sake of analysis of these e model parameters, and to study their dependence on superficial gas velocity, the n means and standard deviations of D_G and K_La at each process rate are calculated and reported in Tables 13 and 14 for all the cases that have been considered in this s study. The averages are calculated for both trials of each run, excluding the parameters obtained by fitting the response of the lowest detector.

Table 13: Average D_G at each Process Rate for the Various Cases Studied

Case l No.	Run	D_G	σ_G
		cm^2/s	cm^2/s
1 .	14.6	5637	1305
	14.7	2801	1010
	14.8	4951	1504
2а з	14.6	6601	835
	14.7	2984	873
	14.8	9795	1770
2b b	14.6	5650	446
	14.7	2330	204
	14.8	8072	1155
3 +	14.6	6094	$\overline{992}$
	14.7	3236	722
	14.8	7684	1047
4 :	14.6	6098	1019
	14.7	3223	952
	14.8	6388	1301
5 i	14.6	4816	473
	14.7	2910	336
	14.8	6861	1392
6;	14.6	5649	1022
	14.7	2908	723
	14.8	6621	1437

Except for Case 1, all other cases show an increase in D_G with gas velocity. Case 1 has H as one of the floating parameters. The fitted values of H are considerably lower than their corresponding values from thermodynamics. This may be the reason for this anomaly. In general, \mathfrak{t} there is a very large spread of D_G about the mean value.

Literature correlations foor D_G (Mangartz an Pilhofer, 1980; Field and Davidson, 1980; Towell and Ackerman, 1972), reported in Table 15, are valid only for superficial

Table 14: Average $K_L a_i a$ at each Process Rate for the Various Cases Studied

Casese No.	Run	$K_L a$	σ_k
		1/s	1/s
1	14.6	1.23	0.81
	14.7	0.91	0.92
	14.8	1.18	0.97
2 2a	14.6	2.7	2.2
	14.7	1.66	1.00
	14.8	4.19	4.04
22b	14.6	0.358	0.26
	14.7	0.195	0.03
<u> </u>	14.8	1.834	0.99
; 3	14.6	1.04	0.79
	14.7	0.44	0.50
	14.8	0.161	0.08
٠ 4	14.6	0.984	0.99
!	14.7	0.349	0.27
	14.8	0.299	0.08
↓ 5	14.6	0.413	0.41
	14.7	0.421	0.14
	14.8	0.494	0.41
(6	14.6	0.370	0.37
	14.7	0.378	0.06
	14.8	0.359	0.12

gas velocities up to 13 cm/ɛ/s, which is lower than the range of velocities of interest to us. Therefore a comparison is made only between Run 14.7, which is at the lowest gas velocity, and the correlation predictions from the literature.

The values of D_G predilicted by the above correlations are much higher than the mean value of $D_G = 2908cmm^2/s$ at U_G of 14 cm/s obtained in this study. The implications of these results are coompletely different from those for the liquid phase results, which suggest that an increase in pressure causes an increase in the liquid dispersion coefficients. The correlation equations suggest that an increase in pressure, which increases the holdup, will lower the magnitude of the gas dispersion coefficients. This occurs because the correlatitions relate the dispersion coefficient to the swarm velocity of the bubbles, which decreases with a decrease in bubble size or an increase in gas holdup. The physical basis for these equations is not completely known. Intuitively it is expected that the pressence of an excess of smaller bubbles should only increase the dispersion of the gas as s more bubbles follow and recirculate along with the liquid

Table 15: Correlations i for Gas Dispersion Coefficient in Bubble Columns

Investigator	Equationn (in SI)	Range of	Prediction of D_G
		Variables	at U_G 14 cm/s
_			(cm^2/s)
Towell and	$D_G = 199.7 D_C^2 U_G$	$0.00854 \le U_G \le 0.13 \text{ m/s}$	5835
Ackerman		$0.0072 \le U_L \le 0.0135 \text{ m/s}$	
		$D_C = 0.406, 1.067 \text{ m}$	
Field and	$D_G = 56.4 \ l \ D_C^{1.33} (\frac{U_G}{\epsilon_G})^{3.56}$	$0.00854 \le U_G \le 0.13 \text{ m/s}$	13435
Davidson		$0.0 \le U_L \le 0.0135 \text{ m/s}$	
		$0.076 \le D_c \le 3.2 \text{ m}$	
Mangartz and	$D_G = 50.0 \ l D_C^{1.50} (\frac{U_G}{f_G})^{3.00}$	$0.015 \le U_G \le 0.13 \text{ m/s}$	11911
Pilhofer		$0.0072 \le U_L \le 0.0135 \text{ m/s}$	
		$0.092 \le D_{\rm c} \le 1.067 \text{ m}$	

phase, while the larger bubbbles move up the column. It must be noted that there is also a large variation in the ppredictions of D_G from correlations. The correlations are shown to have an error of uppto 60 %. Therefore, no conclusion regarding D_G can be made based on the comparison with these correlations.

Unlike the consistent treends for the gas phase dispersion coefficients with superficial gas velocity, $K_L a$ values s shows no particular pattern (as indicated by the results of ANOVA in Tables A.4.3 a and A.4.4 for Cases 1 and 6). In fact the standard deviations indicate the large variriation of $K_L a$ from the mean values. This large spread of the data, and the lack of any trend suggests that the model is quite insensitive to $K_L a$. If one looks at there results from Case 5 and Case 6, which are the most reliable cases, it appears that there is no dependence of $K_L a$ on gas velocity. Here again, the spread is significantly large. Unfortunately, there are no existing literature correlations for $K_L a$ under these conditions of pressure and gas velocity, to make any quantitative comparisons.

5.2.8 Conclusions and FFuture Work

• In general, good fits anre obtained with the one-dimensional model using three floating parameters, $(LD_G, H \text{ and } K_La)$, $(D_G, K_La \text{ and } \epsilon_G)$, $(D_G, K_La \text{ and } \alpha)$ and $(D_G, K_La \text{ and } U_{GG})$. Although the fits are good for these cases, there is a mismatch between the e model predicted H, ϵ_G and α from the first three cases (1, 3 and 4) and corresponding independent measurements. Only for Case 6 are the fitted values of U_G reasonable and relate well to independent experimental informatition.

Since independent esistimates of all model parameters except D_G and $K_L a$ are available, ideally a tytwo floating parameter model should suffice. When such a two parameter moodel, with D_G and $K_L a$ as floating parameters, is used, reasonable fits are possible only when the variation of superficial gas velocity, due to gas consumpticion, is accounted for. A linear decrease in U_G over the entire dispersion height is considered for this. This demonstrates the importance of accounting for the varaying gas velocity in the column, despite the fact that the observed conversisions are quite low. In Equation 13, a decrease in U_G is equivalent to an increase in ϵ_G . This explains the results obtained for Cases 2 and 3. Therefore, by ϵ accounting for the change in ϵ with axial position in the column, the model is able to provide good fits of data when estimates for ϵ from DP measurements are used.

A comparison of the z results between Cases 5 and 6 shows that the D_G values are about the same.

- In the process of using the various approaches for fitting the data, it is shown (in Case 4) that the assumptions for estimating the total argon concentration, given by Equation 7, are good approximations for the system under consideration.
- The values of the parameters obtained by model fits of experimental data are more consistent for the top two to three detector levels for all the cases considered. At the lowever levels the values of D_G start increasing. Based on the dispersion Peclet number, it is seen that the model is most suitable for interpretation of the measurements at the upper levels, where the deviation from plug flow is the smallest.
- A parametric sensitivivity is performed for Case 1, with three floating parameters, D_G , H and $K_L a$, by fiftxing ϵ_G from DP measurements. It is found that the model is most sensitive to HHenry's law constant H, and least sensitive to $K_L a$. Due to this, variations in I $K_L a$ do not affect the model predictions appreciably. It is noted, however, that I $K_L a$ values of zero cannot be used as the data then simply cannot be fitted by thhe model. In addition, based on the results for Case 2 and Case 5, where two flooating parameters are used, considerable model sensitivity to gas holdup and supperficial gas velocity is evident.
- The average values of D_G indicate that D_G increases with superficial gas velocity. Literature correlations predict gas dispersion coefficients much higher in

magnitude than observed in this study. This is thought to be due to the large error involved in the coorrelations.

While there is a consistent increase in D_G with gas velocity, the increase between runs 14.6 and 14.8, that is, for velocities 25.3 cm/s and 36 cm/s appears only marginal considering the large difference in gas velocities (when compared to the differences between runs 14.7 and 14.6), as seen in Table 13. Run 14.8 is at a lower pressure (3.6MHPa) than the other two runs (5.2MPa). This implies that a reduction in pressure causes a lowering of D_G , which for the case of run 14.8 would be expected to 1 be higher at a pressure of 5MPa. However, there is no sufficient experimental 1 information at present to confirm this hypothesis fully.

• While the gas dispersion coefficient seems to show reasonable trends, the volumetric mass transfer coefficient, $K_L a$, shows no pattern. In addition, there are large variations of $K_L a a$ about the mean value.

Overall, the dispersion model is able to match the tracer responses in the column. Judging from all the cases four different floating parameters that are studied, it is clear that an accurate estimate off the average gas holdup and superficial gas velocity is absolutely necessary for goodd fits of data.

With regard to the two maain parameters, a large variation in D_G values is obtained at different axial positions of measurement, which shows an increase of D_G with a decrease in height. The $K_L a a$ values obtained by fitting the data show no consistent patterns and have a very large spread about their means, which basically indicates that the model is insensitive e to $K_L a$. As a result, the effect of U_G on $K_L a$ cannot be properly assessed. This is poerhaps also due to the lack of a physical basis for using the axial dispersion model foor the fluid dynamic situation in the present column.

A model that better capptures the nature of flow in bubble columns and distinguishes between the possibly different bubble sizes is required to consistently predict the characteristics of the tracer responses at all levels in the column. For this purpose, we propose to use the ephenomenological Two Phase Recycle with Cross-Flow Model (TRCFM), which accounts for the movement of different bubble classes (gas) within the column and theirir interaction with the liquid phase. The axial variation of superficial gas velocity annd gas holdup can be suitably incorporated into such a model.

6 Gas - Liquid Mixing and Scale-Up Issues in Bubble Columns

In bubble columns liquidid recirculation is set up in the column in a time-averaged sense (Devanathan et al, 1990; Hills, 1974). The maximum liquid velocity increases with increasing superficial l gas velocity, as the gas holdup increases and the radial holdup profile becomes stateeper (Kumar et al, 1994). Simultaneously, due to increased turbulence there is an increase in the radial and axial turbulent diffusivities (Devanathan, 1991; Degalecesan et al. 1995). The effective or overall axial liquid dispersion coefficient is a result of convective and turbulent mixing. The overall liquid axial dispersion coefficient t can be tentatively approximated by Taylor diffusivity as

$$D_L = \frac{2\overline{U}_{LU}^2 R^{*2}}{D_{zz}} + D_{zz} \tag{22}$$

where \overline{U}_{LU} is the mean upfollow liquid velocity calculated as

$$\overline{U}_{LU} = \frac{\int_0^{R^*} u_{Lz}(r)\epsilon_L(r)rdr}{\int_0^{R^*} \epsilon_L(r)rdr}$$
(23)

As U_G increases, \overline{U}_{LU} , , $D_{\tau\tau}$ and D_{zz} increase. As long as the increase in \overline{U}_{LU}^2 is larger than the increase in $LD_{\tau\tau}$, then both terms in Eqn 22 will increase with increasing gas velocity.

We have experimental I results for liquid velocities and turbulent diffusivities in smaller diameter columns a and lower gas velocities (Degaleesan and Duduković, 1995) velocities. When extrapolatated to larger columns, this yields turbulent diffusivities in the range $D_{rr} \sim 100$ to $1500~cm^2/s$, and $D_{zz} \sim 1500$ to $2000~cm^2/s$. The mean upflow velocities can be obtained uusing the one-dimensional model of Kumar et al. (1994) for liquid velocity. This yields s $\overline{U}_{LU} \sim 23$ to 35 cm/s. The above values are calculated for the existing conditions from U_G 14 cm/s to 36 cm/s, respectively. When substituted in Eqn 22, this yields values of D_L to be 3300 to 6100 cm^2/s . It is noted that the expression for Taylor diffusisivity is only an approximate one for use in bubble columns. In addition the convective a and turbulent parameters are obtained by extrapolation of the existing data for smalldler diameter columns. This gives us reasonable results for D_L under the existing operating conditions.

Using the axial dispersision model (ADM) and fitting the model to experimental data, the axial dispersion 1 coefficients for the gas, D_G , and liquid, D_L , have been evaluated from tracer data 1 under methanol synthesis conditions. For the liquid phase,

the relatively larger values ϵ of dispersion coefficients obtained, when compared to correlation predictions of liquid dispersion coefficients at atmospheric pressure, are qualitatively justified based ϵ on the holdup differences between the two systems. For the gas phase, comparison obf the fitted D_G with literature correlations are counter intuitive and opposite to that of the liquid phase experimental results.

It is clear that none of the existing correlations can be used satisfactorily to predict mixing (dispersion coefficients) in the gas or liquid phase under existing conditions. Since the experiments considered here are only for a single column diameter, and since the effect of pressure can noot be discerned from these experiments, no correlations can result from the present; work. In addition, results of parameter estimation at the various detector levels inadicate the effect of L/D ratio on the model parameters, especially at the lower detector levels of the column.

With regard to scale-up t to industrial size columns, caution must be excerised in extrapolating the correlations of the dispersion coefficients to large diameter columns.

7 Conclusions

The liquid (slurry) and ga;as phase tracer data have been interpreted with the one dimensional axial dispersion 1 model (ADM).

For the case of the liquidd, the dispersion model is able to give good fits for the detectors that do not exhibibit any overshoots. The obtained dispersion coefficients show an increase with gas vivelocity. However, there is a large scatter in the axial dispersion coefficient obtained at different detector levels, which suggests that the model is inadequate to describe the behavior of the liquid within the entire column with a single dispersion coeffficient. For this purpose, we propose to use the Recycle with Cross Flow and Dispersision Model (RCFDM).

For the gas phase tracer analysis, good fits are possible using the ADM at almost all detector levels, but there i is a variation in the estimated parameters with detector level. While the gas phase dispspersion coefficient shows an increase with gas velocity, no particular dependence is seenn with regard to the volumetric mass transfer coefficient. To better represent the flow pattern along with the bubble size distribution and gas - liquid exchange, we proposes to use the Two Phase Recycle with Cross Flow model (TRCFM) in the future.

Comparison with literatuure correlations for the dispersion coefficients shows the lack of any suitable correlations under existing operating conditions. In order to use the correlations for scale-up,, additional experiments at different column diameters,